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# Design and Development of a Magneto-Optic Sensor for Magnetic Field Measurements

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**Abstract:** A magneto-optic sensor is developed using a Terbium Doped Glass (TDG) element as a Faraday rotation sensor and optical fiber as light transmitting and receiving medium. Online LabView based application software is developed to process the sensor output. The system is used to sense the magnetic field of a DC motor field winding in industrial environment. The sensor output is compared with the magnetic flux density variation obtained with a calibrated Hall Magnetic sensor (Gauss Meter). A linear variation of sensor output over wide range of current passing through the field winding is obtained. Further the results show an improved sensitivity of magneto-optic sensor over the Hall sensor. *Copyright* © 2015 IFSA Publishing, S. L.

**Keywords**: Current measurement, Faraday effect, Magnetic field measurement, Magnetooptic devices, Magnetic sensors, Optical imaging, Polarimetry, Stokes parameters.

#### 1. Introduction

Different technologies are used for measuring magnetic fields. Lion [1] described fundamental description of both mechanical and electrical means for sensing magnetic fields. H. R. Everett and J. E. Lenz [2, 3] have described surveys on magnetic sensor technologies. Induction coil magnetometer [4], Fluxgate magnetic sensor [4], Hall magnetic sensor [5] etc are commercially available magnetic field sensors used for wide ranging applications such as measuring rock magnetism, satellite altitude control Super-conducting system etc.[6] interference device [3], Giant magneto resistance (GMR) magnetic field sensor [7-9] etc are some important developments in magnetic technology. The optical sensors are now-a-days gaining importance since they are free from measurement errors due to various effects like stray

electromagnetic effect, capacitive effect, etc., which are observed in many non-optical sensors along with their advantages of electrical isolation, large bandwidth, ease of integration into digital control system, etc. [10–15].

One very important property used in optical sensors is the Faraday rotation of polarized light in a typical solid, liquid or gaseous medium such as glass, quartz, water, and sodium vapor when subjected to a strong magnetic field [16]. Ferrimagnetic materials such as Yttrium Iron Garnet (YIG) (Y<sub>3</sub>Fe<sub>15</sub>O<sub>12</sub>) and ferromagnetic materials such as Iron (Fe), Nickel (Ni) and Gadolinium (Ge) offer highest Faraday rotation but high degree of nonlinearity is also observed in the net Faraday Effect [17]. Large iron garnet crystal is not suitable where sensor linearity is important.

In these materials volumes of equal direction of magnetization (magnetic domains) are formed.

The external magnetic field applied perpendicular to such material cause the domains with magnetization in the direction of the field to grow at the expense of other domains [17]. The net polarization rotation can be approximated using the area ratio between two kinds of domain. The domain size and hence the Faraday rotations depend on the material property as well as sample geometry. Owing to high Faraday rotation these materials are also used in thin film geometry with compensation technique for nonlinearity. The Faraday rotations can be altered by adding mainly rare earth ions. Bismuth (Bi) is commonly added to achieve higher rotations. Jeon et al. showed that Faraday rotation and Verdet constant of Bi-substituted YIG (Bi<sub>1.8</sub>Y<sub>1.2</sub>Fe<sub>5</sub>O<sub>12</sub>) material can be enhanced to 5.28 deg/mm and 2.63×10<sup>-2</sup> deg/Oe cm [18]. Some films are made with complete substitution of Bi in place of Y such as the composition Bi<sub>3</sub>Fe<sub>5</sub>O<sub>12</sub> or BIG which commercially used in laser isolators [19]. In magneto-optic current sensors the principle of Faraday rotation is extensively used by using bulk optical element or optical fiber [20-24] as magnetooptic medium. However, there is a constraint in the use of multi-turn optical fiber due to low Verdet constant and inherent intrinsic birefringence and bend-induced linear birefringence effect [25]. This leads to the use of bulk optics in place of optical fiber [26-30]. Chakraborty et al. [31-35] designed and developed magneto-optic sensors using bulk optical material. The diamagnetic materials like SF-57, SiO<sub>2</sub> and BK7 etc. and paramagnetic materials like Terbium Gallium Garnet (TGG) crystal, Terbium doped glass (TDG), etc. used in bulk optics are still being studied by various workers [26–30]. Gillman has observed large rotations in magneto-optic laser isolators at 632 nm wavelength using TGG and Tbdoped glass (TDG) of Verdet constant -134 and -70 rad/ T.m respectively [26]. Mansell et al. have observed that change in optical path length due to thermal expansion is smaller in Hoya FR-5 (TDG) than TGG [28]. Very recently Li et al. have developed Faraday glasses with large size and high performance [29]. Goldring et al. have utilized magneto-optic element for control of polarization [30]. In the present work the design, development and testing of a prototype TDG based magneto-optic

sensor is discussed. A comparative study with the performance of Hall sensor is also included.

#### 2. Theory and Analysis

Basic schematic diagram of the magneto-optic measurement set up in transmission mode is shown in Fig. 1. Let the beam incident on the magneto-optic element be linearly polarized with azimuth  $\alpha_p$  has intensity  $I_0$ . The Stokes vector of the output beam from the analyzer can then be written as

$$S_{out} = I_0 M_{ana} M_{rot} S_{in}, \qquad (1)$$

where  $M_{rot}$  and  $M_{ana}$  are the Mueller matrices [13, Chapter-33] of Faraday rotator and analyzer respectively, and  $S_{in}$  is the normalized Stokes vector of the input linearly polarized beam which are given by the following equations

$$M_{rot} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos 2\theta & -\sin 2\theta & 0 \\ 0 & \sin 2\theta & \cos 2\theta & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$
 .(2)

$$M_{ana} = \frac{1}{2} \begin{bmatrix} 1 & \cos 2\alpha_{a} & \sin 2\alpha_{a} & 0\\ \cos 2\alpha_{a} & \cos^{2} 2\alpha_{a} & \cos 2\alpha_{a} \sin 2\alpha_{a} & 0\\ \sin 2\alpha_{a} & \cos 2\alpha_{a} \sin 2\alpha_{a} & \sin^{2} 2\alpha_{a} & 0\\ 0 & 0 & 0 & 0 \end{bmatrix}$$
(3)

$$S_{in} = \begin{bmatrix} 1 \\ \cos 2\alpha_p \\ \sin 2\alpha_p \\ 0 \end{bmatrix}, \tag{4}$$

where  $\alpha_a$  is the transmission angle of the analyzer, and  $\theta$  is the Faraday rotation.

Considering linearly horizontally polarized  $(\alpha_p = 0)$  monochromatic input beam and combining the Eqs. (1)–(4), the expression for detected intensity of the out put beam is given by

$$I = \frac{I_0}{2} \left[ 1 + \cos 2(\theta - \alpha_a) \right] \tag{5}$$

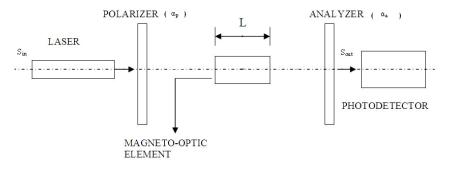


Fig. 1. Basic schematic diagram of the magneto-optic measurement set up.

when a linearly polarized light passes through a magneto optic medium in a direction parallel to the uniform magnetic field, the Faraday rotation is given by the relation [36]

$$\theta = V_{verdet} LB, \qquad (6)$$

where,  $V_{\text{Verdet}}$ , L and B are the wavelength dependent Verdet constant of the TDG element, length of the magneto-optic medium and magnetic flux density of the permanent magnet respectively. In order to improve the response and avoid the effect of intensity variation of incident light, the measurement setup utilizes the dual quadrature polarimetric detection scheme in reflection mode where the intensity of the output beam is measured for  $\alpha_a = +45^\circ$  and  $-45^\circ$  and the ratio of the difference over sum is evaluated [25].

Thus substituting  $\alpha_a = +45^{\circ}$  and  $-45^{\circ}$  in Eq. (5) and taking the ratio of difference over sum, the expression for power ratio (PR) is obtained as

$$PR = \frac{I_{+45} - I_{-45}}{I_{+45} + I_{-45}} = \sin(2\theta)$$
 (7)

In a double pass mode the expression for power ratio will be

$$PR = \sin(4\theta) \tag{8}$$

#### 3. Experiment and Discussions

The dual quadrature polarimetric scheme [25] is shown in Fig. 2.

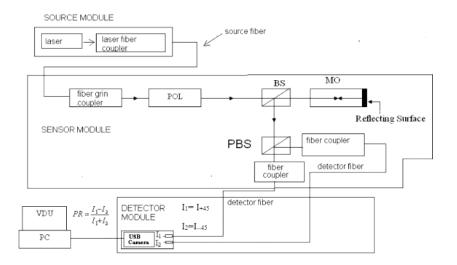
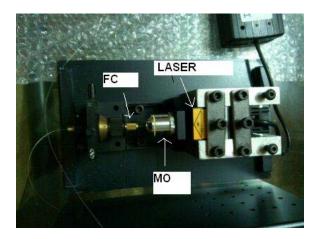


Fig. 2. The dual quadrature polarimetric scheme of the magneto-optic sensor.

source module the light from a 5 mW DPSS (diode pumped solid state) laser (wavelength 0.5435 um) is coupled to a multimode optical fiber (source fiber) by a laser fiber coupler with the help of microscope objective. The other end of the fiber is coupled to the fiber GRIN lens coupler of the sensor module for producing collimated laser beam. The beam passes through a polarizer (POL) and then through a beam splitter (BS) to the magnetooptic (MO) TDG rod of length 30 mm and diameter 10 mm. The reflected beam from the mirrored surface of the magneto-optic (MO) element then passes through the polarizing beam splitter (PBS). The two light beams with orthogonal polarization are then coupled to two multimode optical fibers (detector fibers) with the help of two microscope objective based fiber laser couplers. The other ends of the detector fibers mounted in fiber chucks are focused on to the CCD sensor surface in an USB based camera detector. The advancement in digital image processing, availability of high speed computer with high memory size, simple image grabbing techniques and simplified image storage, manipulation and display enables the use of CCD camera as a light detector in place of single photo detector [37-39] The light intensities  $I_1$  and  $I_2$  are obtained as average gray values of the spot images produced by the camera detector and the power ratio (PR) is measured with the help of the relation  $PR = (I_1 - I_2)/(I_1 + I_2)$ . Figs. 3-5 shows the photographs of prototype source, sensor and camera detector units mounted on nonmagnetic base plates (PMMA materials). For online measurement of PR, application software is developed in LabView utilizing IMAQ vision tools. The sensor system is first tested in the laboratory by inserting the TDG cylinder inside the core of a multiple layer solenoid coil (copper) of length 95 mm and with 2600 turns.

Fig. 6 (a) shows the variation of PR with the variation of current (DC) through the solenoid coil. The variation of magnetic flux density with the variation of coil current is also measure by the axial Hall probe of a calibrated Gauss meter and a maximum linear change of around 35 mT is obtained when the current is changed from 0 to 2 Amp. Fig. 6 (b) shows the image of the fiber tips acquired

in LabView for calculating the light intensities  $I_1$  and  $I_2$  from the average gray values within the respective region of interest.



**Fig. 3.** Source module mounted on the base plate.

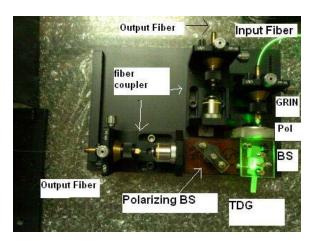


Fig. 4. Sensor module mounted on the base plate.

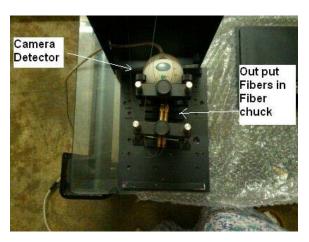
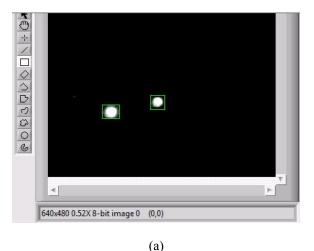
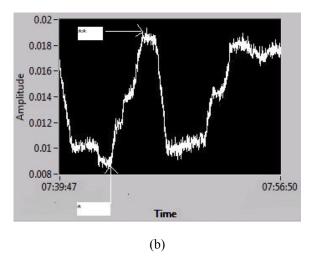


Fig. 5. Detector module mounted on the base plate.

Fig. 7 shows the experimental setup for sensing the DC motor field winding. The magneto-optic TDG element is located at a distance 10 mm from DC motor field winding and input current to the field

winding is increased slowly from 1 ampere to 7 ampere in steps of 0.5 Ampere. The experiment is repeated for increasing and decreasing current. For comparison with the data obtained from magneto-optic setup an axial Hall probe of a calibrated gauss meter is also used to sense the magnetic flux density variation as the current changes. The resulting PR is obtained by using Eq. 6 and 8 and considering TDG element of length L=30 mm and verdet constant  $V_{verdet} = 101.81 \text{ rad/T.m}$  at  $\lambda$ =0.5435  $\mu$ m [31]. In the measurement process the residual magnetism of 0.01K gauss in the field winding is obtained.





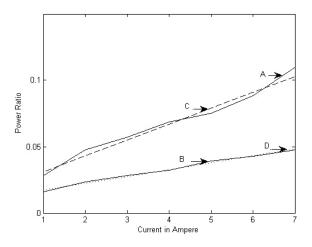
**Fig. 6.** Testing with a solenoid current of 2 Amps in laboratory. (a) Shows the variation of PR with time when the solenoid coil current is changed from 0 -2 Amp. At point \* the coil current is practically zero and at point \*\* the coil current is 2 Amps (b) Shows the image of the detector fiber tips surrounded by region of interest for calculating the average gray value.

Curve A of Fig. 8 shows the average of ten measurements using magneto-optic sensor element and for both ascending and descending order of current. Curve-C shows the linear fit of Curve-A with sum of the square of errors (SSE) =9.662e-005 and

R-square=0.9778. Curve-B of Fig. 8 shows the average of five measurements using Hall sensor probe and for both ascending and descending order of current. Curve-D shows the linear fit of Curve-B with sum of the square of errors (SSE)= 5.325e-006 and R-square=0.9929.



**Fig. 7.** Experimental setup for detecting the change in coil current in the field winding of a DC motor.



**Fig. 8.** Variation of Power Ratio (PR) with the variation of coil current (of field windings). Curve-A (—) shows the variation obtained with magneto-optic TDG element and Curve-C (——) shows the linear fit of Curve-A. Curve-B (—) shows the variation obtained with a Hall magnetic sensor and Curve-D (….) shows the linear fit of Curve-B.

Both the curves show a fairly linear variation over a wide range of current passing through the field winding. Interestingly we note that the sensitivity of the magneto-optic sensor is practically more than double the sensitivity of Hall probe sensor since the slopes of the Curves C and D are 40° and 17° respectively. The difference between Curve-A and B can be explained by the improper polarizer setting ("Pol" in Fig. 4) of input polarizer with transmission

angle  $\alpha_p$  which we have considered as linearly horizontally polarized i.e.  $\alpha_p = 0^\circ$ . With the finite value of  $\alpha_p$  the expression for Power Ratio (PR) will be changed to PR=  $\sin{(4\theta+2\alpha_p)}$ . The magneto-optic sensor developed can further be miniaturized by replacing the bulky fiber coupler (Fig. 4) with a lens and molded plastic packaging.

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